

## RELATION BETWEEN THE DOA MATRIX METHOD AND THE ESPRIT METHOD

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### ABSTRACT

There have been many algorithms dealing with the estimation of the direction-of-arrival (DOA) for multiple narrowband sources. Recently the most well-known of these is undoubtedly the ESPRIT method. Moreover, a new 2-D eigenstructure approach called the DOA matrix method has been proposed. Unlike the ESPRIT method, it estimates the 2-D directions from both nonzero eigenvalues and corresponding eigenvectors of a DOA matrix. In this paper the relation between them is theoretically analyzed. It is shown that the ESPRIT method may be regarded as a special case of the DOA matrix method, and that the DOA matrix method may give more generalized and more perfect results.

### I. Introduction

Estimation of the direction of arrival (DOA) from noisy sensor array data has attracted tremendous research attention for several decades because of its application in radar, sonar, seismic, and radio signal processing. There have been a variety of techniques and algorithms proposed for dealing with this issue. The sensor arrays may not be uniformly spaced and linear. But, with the uniformly spaced linear sensor array systems, a very important array system in practice, we can save a lot of computations and storages. Here, we discuss the linear sensor array system.

In recent years, there has been a growing interest in eigenstructure based methods. These methods, pioneered by Pisarenko [1], Schmidt [2] and Kumaresan [3], are known to yield high resolution and asymptotically unbiased estimates. Furthermore, Paulraj, et. al., [4,5,6] have proposed a subspace rotation approach called ESPRIT. Like the MUSIC method [2], it exploits correctly the underlying signal model to generate asymptotically unbiased estimates. But it estimates the DOA by using the eigenvalues of a matrix pair rather than the eigenvectors of the covariance matrix,

which are used by most of the eigenstructure methods including the MUSIC method. Without the one-dimensional (1-D) spectral peak search, the ESPRIT method has better performance and less computations than previous ones. But these mentioned above are 1-D methods. When applying them to linear sensor arrays, we have to assume that all radiating sources are located in the same plane. On the other hand, most available two-dimensional (2-D) methods need a large rectangular plane sensor array and have to perform 2-D spectral peak searches with large computations.

Recently a new 2-D eigenstructure approach called the DOA matrix method was proposed in [7,8] and extended to the 3-D case [9] with good performance in resolution and computation. Unlike the ESPRIT method, the DOA matrix method estimates the 2-D angles of arrival by simultaneously using the nonzero eigenvalues and corresponding eigenvectors of a DOA matrix. It seems that there is some relation between these two methods.

In this paper the ESPRIT method and the DOA matrix method are briefly introduced. The relation between them is theoretically analyzed. It is pointed out that the ESPRIT method is just a special case of the DOA matrix method, and the DOA matrix method gives more generalized and more perfect results. Both should have the same performance for 1-D cases. But the DOA matrix method can be used in the 2-D case. Furthermore, it even can be extended to solve the 3-D problem, as we did in [9], while the ESPRIT method can only be used in the 1-D case.

### II. Problem Formulation

Consider an array system consisting of two uniformly spaced linear subarrays of  $p$  sensors spaced  $D$  apart as shown in Fig. 1. The two subarrays, the subarray  $X_a$  and the subarray  $Y_a$ , are parallel with subarray  $X_a$  lying on the  $X$  axis and starting at the origin;  $d$  is the distance between them as measured by the  $Y$  axis intercept of subarray  $Y_a$ . Assume that plane waves with known center frequency  $\omega_0$  emitted by  $K$  narrowband sources impinge on this array system and that the DOA of the

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sources are  $\{\theta_1, \theta_2, \dots, \theta_K\}$ , where  $\theta_k = (\alpha_k, \beta_k)$ ,  $\alpha_k$  is the DOA of the  $k$ -th source relative to the  $X$  axis and  $\beta_k$  relative to the  $Y$  axis, as shown in Fig. 1. The sampled vectors of subarrays  $X_a$  and  $Y_a$  can be written as [7]

$$X(t) = As(t) + N_x(t) \quad (1)$$

$$Y(t) = A\Phi s(t) + N_y(t) \quad (2)$$

where

$$X(t) = [x_1(t), x_2(t), \dots, x_p(t)]^T \quad (3)$$

$$Y(t) = [y_1(t), y_2(t), \dots, y_p(t)]^T \quad (4)$$

$$s(t) = [s_1(t), s_2(t), \dots, s_K(t)]^T \quad (5)$$

$$N_x(t) = [n_{x1}(t), n_{x2}(t), \dots, n_{xp}(t)]^T \quad (6)$$

$$N_y(t) = [n_{y1}(t), n_{y2}(t), \dots, n_{yp}(t)]^T \quad (7)$$

$A$  is a  $p \times K$  Vandermonde matrix

$$A = [a(\alpha_1), a(\alpha_2), \dots, a(\alpha_K)] \quad (8)$$

with  $k$ -th column

$$a(\alpha_k) = [1, \dots, \exp(j\omega_0(p-1)\tau_{Dk})]^T \quad (9)$$

$\tau_{Dk}$  is the inter-element path delay of the plane wave from the  $k$ -th source,  $\tau_{Dk} = (D/c)\cos\alpha_k$ , where  $c$  is the wave propagation velocity.  $\Phi$  is a  $K \times K$  diagonal matrix

$$\begin{aligned} \Phi &= \text{diag}[\exp(j\omega_0\tau_{d1}), \dots, \exp(j\omega_0\tau_{dK})] \\ &= \text{diag}[\phi(\beta_1), \dots, \phi(\beta_K)] \end{aligned} \quad (10)$$

$\tau_{dk}$  is the path delay of the plane wave of the  $k$ -th source between the two subarrays,  $\tau_{dk} = (d/c)\cos\beta_k$ .

Note that the  $k$ -th column of the matrix  $A$ ,  $a(\alpha_k)$ , and the  $k$ -th element,  $\phi(\beta_k)$ , of the diagonal matrix  $\Phi$  are associated with the DOA, i.e.  $(\alpha_k, \beta_k)$ , of the  $k$ -th source. They are referred to as the signal vector and signal element, respectively.

The auto-covariance matrix of  $X(t)$  is given by

$$R_{xx} = E[X(t)X(t)^H] = ASA^H + \sigma^2 I \quad (11)$$

where  $^H$  denotes the conjugate transpose and  $E$  the expectation operator,  $S = E[s(t)s(t)^H]$  is the covariance matrix of source signals,  $I$  is the identity matrix, and  $\sigma^2$  is the variance of the additive noise. The cross-covariance matrices of  $Y(t)$  and  $X(t)$  are given by

$$\begin{aligned} R_{yx} &= E[Y(t)X(t)^H] \\ &= A\Phi SA^H + E[A\Phi s(t)N_x(t)^H] + E[N_y(t)s(t)^H A^H] \\ &\quad + E[N_y(t)N_x(t)^H] \end{aligned} \quad (12)$$

and

$$R_{xy} = E[X(t)Y(t)^H] = R_{yx}^H \quad (13)$$

We assume that the additive noises are uncorrelated with signals and with each other, in which case Eqs. (12) and (13) can be written

as

$$R_{yx} = A\Phi SA^H \quad (14)$$

$$R_{xy} = R_{yx}^H = AS\Phi^H A^H \quad (15)$$

### III. The ESPRIT Method

In Eq. (11), the auto-covariance matrix  $R_{xx}$  has two terms. The one associated with source signals is denoted by

$$R_{xx0} = R_{xx} - \sigma^2 I = ASA^H \quad (16)$$

Obviously, the rank of  $R_{xx0}$  is equal to  $K$ , the number of sources if there are no fully correlated (coherent) sources.

The subspace rotation method (ESPRIT) for DOA estimation relies on determining  $\Phi$  from the estimated covariance matrices  $R_{xx0}$  and  $R_{xy}$ . It is shown in [4] if  $S$  is nonsingular, the  $K$  nonzero generalized eigenvalues of the matrix pair  $\{R_{xx0}, R_{xy}\}$  are equal to  $\phi(\beta_k)$ . Thus, if  $\{\mu_i\}$ ,  $i=1,2,\dots,K$ , are the nonzero eigenvalues of the matrix pair  $\{R_{xx0}, R_{xy}\}$ , the estimated angles  $\{\beta_i\}$ ,  $i=1,2,\dots,K$ , are given by

$$\beta_i = \cos^{-1} \left\{ \frac{c}{\Omega_0 d} \arg(\mu_i) \right\} \quad (17)$$

The ESPRIT method is a 1-D method, it merely estimates the 1-D angles,  $\{\beta_i\}$ ,  $i=1,2,\dots,K$ , from the generalized eigenvalues of the matrix pair  $\{R_{xx0}, R_{xy}\}$ .

### IV. The DOA Matrix Method

In [7], a  $p \times p$  matrix  $R$  referred to as the DOA matrix is defined

$$R = R_{yx} R_{xx0}^* \quad (18)$$

where  $R_{xx0}^*$  denotes the pseudoinverse of  $R_{xx0}$  and is constructed by the nonzero eigenvalues,  $\sigma_i$ , and corresponding eigenvectors,  $v_i$ , of  $R_{xx0}$  via

$$R_{xx0}^* = \sum_{i=1}^K \sigma_i^{-1} v_i v_i^H \quad (19)$$

It is shown in [7] that if  $A$  and  $S$  are nonsingular, the DOA matrix  $R$  has its  $K$  nonzero eigenvalues equal to the  $K$  diagonal elements of  $\Phi$  and corresponding eigenvectors equal to the  $K$  column vectors of matrix  $A$ , i.e.

$$RA = A\Phi \quad (20)$$

Therefore, if  $\{\epsilon_i\}_{i=1}^K$  and  $\{u_i\}_{i=1}^K$  are the  $K$  nonzero eigenvalues and corresponding eigenvectors of the DOA matrix  $R$ , the estimated  $\alpha_i$  and  $\beta_i$  can be given by

$$\alpha_i = \cos^{-1} \left\{ \frac{c}{\Omega_0 D(p-1)} \sum_{n=2}^p \frac{1}{n-1} \arg \left[ \frac{u_i(n)}{u_i(1)} \right] \right\} \quad (21)$$

and

$$\beta_i = \cos^{-1} \left\{ \frac{c}{\Omega_0 d} \arg(\epsilon_i) \right\} \quad (22)$$

Note the DOA matrix method estimates both angles,  $\alpha_k$  and  $\beta_k$ , simultaneously. Clearly, in the special case that all the sources are located in the same plane, it can also be used by simply calculating the  $\alpha_k$  or  $\beta_k$ , as with the ESPRIT method.

#### V. Relation Between the Two Methods

Let  $\{\mu_i\}_{i=1}^K$  and  $\{q_i\}_{i=1}^K$  be the nonzero generalized eigenvalues and corresponding eigenvectors of the matrix pair  $\{R_{x \times 0}, R_{xy}\}$  in the ESPRIT method. We have

$$R_{x \times 0} q_i = \mu_i R_{xy} q_i \quad i=1, 2, \dots, K \quad (23)$$

Substitute  $R_{xy}$  with  $R_{yx}$  according to (15)

$$R_{x \times 0} q_i = \mu_i R_{yx}^H q_i \quad (24)$$

Take the complex conjugate transpose of both sides of (24)

$$q_i^H R_{x \times 0} = \mu_i^* q_i^H R_{yx} \quad (25)$$

where \* denotes the complex conjugate. Then

$$q_i^H R_{x \times 0} R_{x \times 0}^* = \mu_i^* q_i^H R_{xy} R_{x \times 0}^* \\ q_i^H \left[ \sum_{i=1}^K \sigma_i v_i v_i^H \right] \left[ \sum_{i=1}^K \sigma_i^{-1} v_i v_i^H \right] = \mu_i^* q_i^H R_{yx} R_{x \times 0}^* \\ q_i^H \left[ \sum_{i=1}^K v_i v_i^H \right] = \mu_i^* q_i^H R_{yx} R_{x \times 0}^* \quad (26)$$

According to the orthogonality between the signal subspace and the noise subspace,

$$q_i^H \perp [v_{K+1}, v_{K+2}, \dots, v_p]$$

and

$$q_i^H v_k = 0 \quad k=K+1, K+2, \dots, p \quad (27)$$

Thus, Eq. (26) can be rewritten as

$$q_i^H \left[ \sum_{i=1}^K v_i v_i^H + \sum_{i=K+1}^p v_i v_i^H \right] = \mu_i^* q_i^H R_{yx} R_{x \times 0}^* \quad (28)$$

Here

$$\left[ \sum_{i=1}^K v_i v_i^H + \sum_{i=K+1}^p v_i v_i^H \right] = I \quad (29)$$

Therefore, we have

$$q_i^H R_{yx} R_{x \times 0}^* = (\mu_i^*)^{-1} q_i^H \quad (30)$$

Note that  $R_{yx} R_{x \times 0}^*$  is just the DOA matrix  $R$  in (18)

$$q_i^H R = (\mu_i^*)^{-1} q_i^H \quad (31)$$

We know that the nonzero eigenvalues of the DOA matrix have a unit magnitude, i.e.  $|\mu_i| = 1$ . Therefore, Eq. (31) can be simplified to

$$q_i^H R = \mu_i q_i^H \quad (32)$$

where  $q_i^H$  is the left eigenvector of  $R$  rather than the right eigenvector which is used by the DOA matrix method. Furthermore,  $q_i^H$  is not the signal vector in (10). Eq. (32) implies that the eigenvalues  $\{\mu_i\}_{i=1}^K$  of the matrix pair in (23) for the ESPRIT method are also the eigenvalues of the DOA matrix  $R$  for the DOA matrix method. In the ESPRIT method, only the eigenvalues are used so that almost half of the information is lost. But in the DOA matrix method, both eigenvalues and eigenvectors are fully utilized to estimate the 2-D bearing angles. From these arguments we know that the ESPRIT method is just a special case of the DOA matrix method, and the DOA matrix method is the generalized ESPRIT method. Therefore, both should have the same performance for the 1-D case. But the DOA matrix method can be used in the 2-D case. Furthermore, it can be even extended to solve the 3-D problem, as we did in [9], while the ESPRIT method can only be used in the 1-D case. Moreover, as both the covariance matrices  $R_{x \times 0}$  and  $R_{xy}$  are not of full rank, it causes trouble to decompose the generalized eigenvalues of the matrix pair in the ESPRIT method, which has been discussed in [10]. But in the DOA matrix method we have no such problem.

It is interesting to note that in almost all of the eigenstructure methods the eigen-decomposition merely separates the signal subspace from the noise subspace, but in the DOA matrix method it gives the original basis of the signal subspace directly. Without requiring any spectral peak searches, the 2-D DOA matrix method is highly efficient in computation. Thus it gives more generalized and more perfect results.

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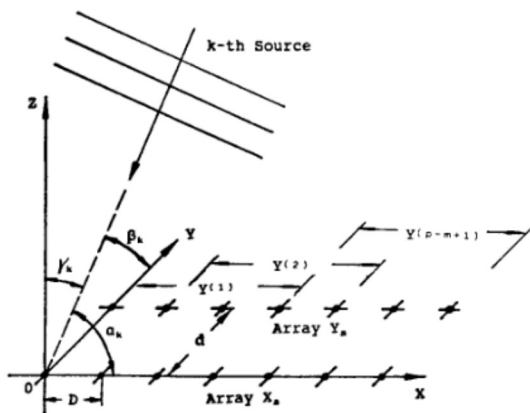


Fig. 1. The geometry of the array system