

B. Dziurla  
 Instytut Informatyki  
 Politechnika Gdanska  
 ul. Majakowskiego 11/12  
 8-952 Gdansk, Poland  
 and  
 Applied Mathematics Program  
 University of Maryland  
 College Park, MD 20742 USA

R.W. Newcomb\*  
 Microsystems Laboratory and  
 Electrical Engineering Department  
 University of Maryland  
 College Park, MD 20742 USA  
 Phone: (301) 454-6869

Abstract:

An example, the first degree Hazony filter section, is used to show that semistate equations become regularized by parasitics in a manner different than postulated in the literature.

I. Introduction

The inclusion of parasitics unnecessarily complicates the design process for some systems, as for example those containing hysteresis. However, in some such cases the practical operation of the finally designed system depends critically upon the presence and effects of parasitics, as for example in determining single-valued hysteresis systems. In the case of semistate described systems, available literature implies that the parasitics lead to a nonsingular coefficient matrix on the semistate derivative (see for example [1, p.112]). Here we show by example that such need not be the case.

II. Canonical Equation Framework

Given a semistate describable system it has the canonical description [2]

$$\begin{aligned} \mathcal{A}\dot{x} + \mathcal{B}(x,t) &= \mathcal{D}u \\ y &= \mathcal{H}x \end{aligned} \quad (1a)$$

where the  $\mathcal{A}$ ,  $\mathcal{B}$ , and  $\mathcal{H}$  are constant matrices,  $\mathcal{A}$  being  $n \times n$  for  $n$ -vector semistates  $x$ ;  $u$  and  $y$  are inputs and outputs, respectively, and  $\dot{\phantom{x}}$  denotes differentiation in time  $t$ .

To tie in with previously published material [1], [3], we bring the coefficient of  $\dot{x}$  to be the identity on the nonsingular part of  $\mathcal{A}$  and zero on the singular portion, as others have done in the linear case [4, p.33][5]. That is, by an equivalence transformation, of nonsingular matrices  $P$  and  $Q$ , we write

$$P\mathcal{A}Q = I_m + 0_{n-m} \quad (2a)$$

$$Q^{-1}x = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}, \quad \begin{array}{l} \text{dimension } x_1 = m \\ \text{dimension } x_2 = n-m \end{array} \quad (2b)$$

where  $I_m$  is the  $m \times m$  identity,  $m = \text{rank } \mathcal{A}$ ,  $+$  denotes the direct sum,  $0_{n-m}$  is the  $(n-m) \times (n-m)$  zero matrix. Then (1) is rewritten as

$$\dot{x}_1 + \mathcal{B}_1(x_1, x_2, t) = \mathcal{D}_1u \quad (3a)$$

$$\mathcal{B}_2(x_1, x_2, t) = \mathcal{D}_2u \quad (3b)$$

$$y = \mathcal{H}_1x_1 + \mathcal{H}_2x_2 \quad (3c)$$

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In [3, p.1112] it is postulated that parasitics add to the left of (3b) the term  $\epsilon \dot{x}_2$  where  $\epsilon$  is a small parameter. In the following examples we show that such need not be the case.

III. An Example

To show the effect on the form of the equations it is initially sufficient to consider a linear example, for which we choose the loaded degree-one Hazony section of Fig. 1a), this section playing an important role in active filter synthesis [6, p. 158]. For the circuit graph of Fig. 1b) we choose the tree as branches 1&2 in which case the technique of [2] gives the canonical semistate equations as ( $G = 1/R$ ,  $g = \text{gyration conductance}$ )

$$\begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ C & -C & 0 \end{bmatrix} \dot{x} + \begin{bmatrix} 0 & g & 1 \\ -g & G & -1 \\ 0 & 0 & -1 \end{bmatrix} x = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} u \quad (4a)$$

$$y = [0, 1, 0]x \quad (4b)$$

where

$$u = i_s, \quad y = v_2, \quad x = \begin{bmatrix} v_1 \\ v_2 \\ i_3 \end{bmatrix} \quad (5)$$

Using

$$P = \begin{bmatrix} 0 & 0 & 1/C \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{bmatrix}, \quad Q = \begin{bmatrix} 1 & 1 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad (6a)$$

$$Q^{-1}x = \begin{bmatrix} v_1 - v_2 \\ v_2 \\ i_3 \end{bmatrix}; \quad x_1 = v_1 - v_2, \quad x_2 = \begin{bmatrix} v_2 \\ i_3 \end{bmatrix} \quad (6b)$$

gives for (3)

$$\dot{x}_1 + \begin{bmatrix} 0 & 0 & -1/C \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = [0]u \quad (7a)$$

$$\begin{bmatrix} -g & G-g & -1 \\ 0 & g & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} u \quad (7b)$$

$$y = [0]x_1 + [1, 0]x_2$$

It is observed that if  $G = 0$ , that is, an open circuit load, then  $x_2$  has a singular coefficient matrix in (7b) and, hence, can not be eliminated from (7a) to obtain state-variable equations.

Although the characterization of parasitics for

even this simple system is a difficult problem in itself, let us assume that the device is to be realized by a fully integrated circuit and operated at low frequencies. Then inductive parasitics can be ignored since capacitive and resistive effects will dominate, leading to the circuit of Fig. 2. In Fig. 2,  $C_1$  &  $R_1$  are Norton equivalent parasitic input capacitance and resistance for the gyrator,  $R_2$  is the similar output resistance,  $C_2$  combines the similar output capacitance with the parasitic capacitance of the load while  $R_3$  is the parasitic resistance of the capacitor C. The same graph may be used as for Fig. 1 resulting in

$$\begin{bmatrix} C_1 & 0 & 0 \\ 0 & C_2 & 0 \\ C & -C & 0 \end{bmatrix} \dot{x} + \begin{bmatrix} G_1 & g & 1 \\ -g & G + G_2 & -1 \\ G_3 & -G_3 & -1 \end{bmatrix} x = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix} u \quad (8a)$$

$$y = [0, 1, 0]x \quad (8b)$$

where (5) still holds to define  $u$ ,  $x$ , and  $y$ . The use of  $P$  and  $Q$  of (6a) brings these to

$$\begin{bmatrix} 1 & 0 & 0 \\ 0 & C_2 & 0 \\ C_1 & 0 & 0 \end{bmatrix} \begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ 0 \end{bmatrix} + \begin{bmatrix} G_3/C & 0 & -1/C \\ -g & G - g + G_2 & -1 \\ G_1 & g + G_1 & 1 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ 0 \end{bmatrix} =$$

$$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} u \quad (9a)$$

$$y = [0, 1, 0,] \begin{bmatrix} x_1 \\ x_2 \\ 0 \end{bmatrix} \quad (9b)$$

It is observed that (9a) is not of the form of (7a, b) with  $Cx_2$  added to the left of (7b). Indeed the  $A$  matrix in (9a) is of a different form and too small a rank. Nevertheless, the third component of  $x$ ,  $x_3$ , can be eliminated in (9) to reduce (9) to state-variable equations, even in the case,  $G = 0$ , in which no such reduction is possible when parasitics are not considered. That is, the parasitics considered in Fig. 2 are adequate to regularize the system but their insertion does not regularize the equations in the manner postulated in the literature.

#### IV. Discussion

The example shown here illustrates that parasitics come into play in the semistate equations in a manner not previously foreseen. Nevertheless semistate theory readily handles parasitics. handles parasitics. Beyond that, semistate theory shows how parasitics can be inserted in design to guarantee unique solutions; this is by inserting parasitics until a  $B_2$  is obtained at (3b) that can be

uniquely solved for  $x_2$ . As such, we are led to answer the questions of regularization for RLC circuits raised by Smale [7, p. 209] since any of his circuits can be placed in the framework of the canonical semistate equations of (1).

#### References

- [1]. S. S. Sastry and C. A. Desoer, "Jump Behavior of Circuits and Systems", IEEE Transactions on Circuits and Systems, Vol. CAS-28, No. 12, December 1981, pp. 1108-1124.
- [2]. R. W. Newcomb, "The Semistate Description of Nonlinear and Time-Variable Circuits", IEEE Transactions on Circuits and Systems, Vol. CAS-28, No. 1, January 1981, pp. 62-71.
- [3]. L. O. Chua and G. R. Alexander, "The Effects of Parasitic Reactances on Nonlinear Networks", IEEE Transactions on Circuit Theory, Vol. CT-18, No. 5, September 1971, pp. 520-532.
- [4]. S. L. Campbell, "Singular Systems of differential Equations, II", Pitman Books, Ltd., 1982.
- [5]. L. R. Petzold and C. W. Gear, "ODE Methods for the Solution of Differential/Algebraic Systems", Sandia Report SAND 82-8051, October 1982.
- [6]. R. W. Newcomb, "Active Integrated Circuit Synthesis", Prentice-Hall, Englewood Cliffs, NJ 1968.
- [7]. S. Smale, "On the Mathematical Foundations of Electrical Circuit Theory", Journal of Differential Geometry, Vol. 7, Nos. 1 & 2, September 1972, pp. 193-210.

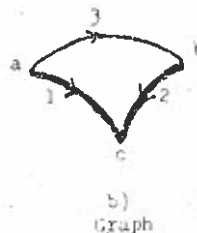
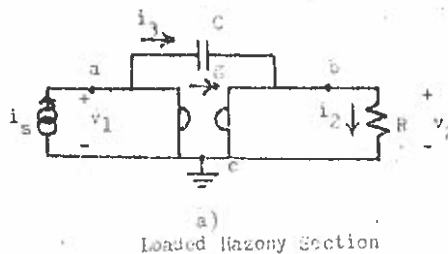


Figure 1

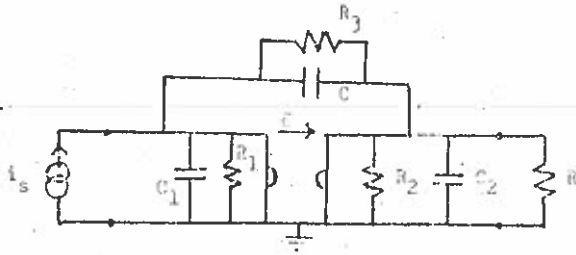


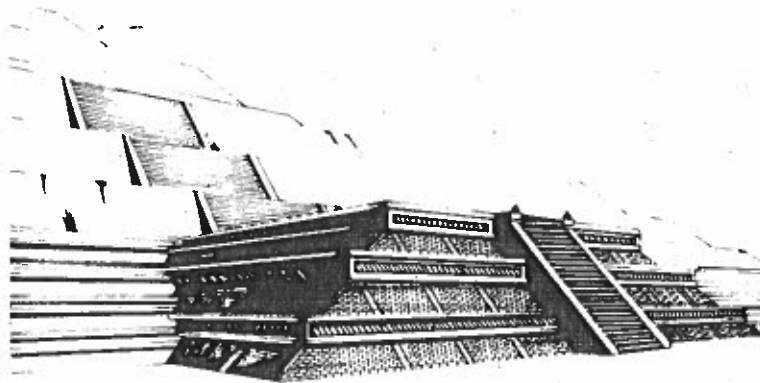
Figure 2  
 Hazy Section with  
 Relevant Parasitics

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CHOLULA PYRAMID

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