

N. C. Debnath
 East Carolina University
 R. C. Ajmera
 East Carolina University
 R. W. Newcomb
 University of Maryland

in Wiley, New
 Springer Verlag,
 Surface Wave
 IEEE Trans.
 No. 11, pp. 856-
 Reversive Inter-
 s. on Micro-
 p. 458-471, July
 y Lines with
 owave Theory
 v. 1969.
 of a SAW
 Proc., pp. 293-

SCATTERING MATRIX FOR N-SECTION
 IN-LINE MODEL OF SAW TRANSDUCERS

ABSTRACT

The scattering matrix of N-equal sections of an in-line model for a SAW device is calculated with the aid of chain and admittance matrices.

1. INTRODUCTION

In the past decade, due to successful applications to the design of different kinds of filters and delay lines, considerable attention has been given to the phenomena of surface acoustic waves (SAW) created on piezoelectric substrates. It has been shown that a circuit model representation can be the key for the development and design of SAW devices. Especially, a frequency domain equivalent circuit model representation [1,2,3] of a transducer has been found to be very successful [2] in characterizing the excitation and detection of surface waves. In these references the transfer characteristic and input admittance at the electric port of the interdigital transducer have been determined with the aid of chain and admittance matrices. Smith, et. al [3], and Gerard [4] have calculated reflection coefficients, but only at the synchronous frequency. Recently, Hribsek and Newcomb [5] have outlined the formalism for obtaining the scattering and transfer scattering matrices of a SAW device section as a function of frequency. While Debnath, Ajmera and Newcomb [6] have calculated the scattering matrix of N-equal sections of a crossed-field model for a SAW section.

The SAW devices offer special interest to system theorists since they rely upon piezoelectrically generated mechanical waves in realizing practical devices for which the theory is anything but complete. And, of course, critical to the operation are the input and output transducers which can be considered as 3-ports having one electrical port and two mechanical (acoustic) ports (one for forward and one for backward acoustic wave motion). These transducers can be characterized by what have been called "in-line" and "crossed-field" models for the basic sections. This paper is concerned with the scattering matrix for

N sections of a SAW transducer described by the in-line model.

The primary objective is to develop a general mathematical technique by using the known admittance matrix of a 3-port SAW transducer section. The starting point is the equivalent circuit of a basic one electrode section of the SAW transducer. This is then used to obtain the 3-port admittance matrix of a full (interdigital) transducer by first finding the matrix of one interdigital period and then applying a cascading formalism. The 3-port admittance matrix for an entire N-section transducer is found by connecting the N sections in cascade acoustically and in parallel electrically. Using these chain and admittance matrix results, the scattering matrix of N equal basic sections expressed in terms of the in-line model is determined.

2. ADMITTANCE MATRIX FOR A THREE-PORT N-SECTION NETWORK

The details of this section repeat those of reference [6] but are recorded here for completeness of the treatment.

We consider a 3-port section with two identical mechanical ports (of variables

$F = \begin{bmatrix} F_1 \\ F_2 \end{bmatrix}$ - force applied,
 $V = \begin{bmatrix} v_1 \\ v_2 \end{bmatrix}$ - velocity of material) and
 one electrical port (of variables $V_3 =$
 voltage, $I_3 =$ current applied to the
 interdigital electrodes), as shown in
 Fig. 1.

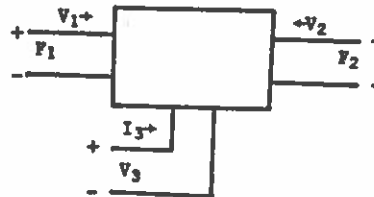


Figure 1. Block-diagram of a three-port Model

Then

$$\begin{bmatrix} v_1 \\ v_2 \\ I_3 \end{bmatrix} = \begin{bmatrix} y_{11} & y_{12} & y_{13} \\ y_{12} & y_{11} & y_{13} \\ y_{13} & y_{13} & y_{33} \end{bmatrix} \begin{bmatrix} F_1 \\ F_2 \\ V_3 \end{bmatrix} \quad (1)$$

is taken as the admittance matrix y where $y_{22} = y_{11}$ and $y_{23} = y_{13}$ by symmetry. y is symmetric by assumed reciprocity.

We now connect N of these sections in cascade at the mechanical ports and in parallel, with every other one reversed, at the electrical ports as shown in Fig. 2.

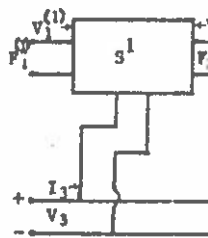


Figure 2. Cascade

Thus for the i^{th} se

$$\begin{bmatrix} v_1^{(i)} \\ v_2^{(i)} \\ I_3^{(i)} \end{bmatrix}$$

while the interconn

where

$$F_1^{(1)} = F_{in} =$$

$$F_2^{(N)} = F_{out}$$

Also for electrical

In order to find th
 reverse - chain mat

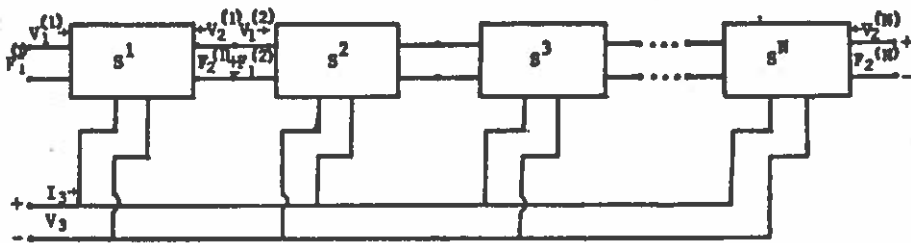


Figure 2. Cascaded one-electrode sections for a SAW transducer.

Thus for the i^{th} section we have

$$\begin{bmatrix} V_1^{(i)} \\ V_2^{(i)} \\ I_3^{(i)} \end{bmatrix} = \begin{bmatrix} y_{11}^{(i)} & y_{12}^{(i)} & y_{13}^{(i)} \\ y_{12}^{(i)} & y_{11}^{(i)} & y_{13}^{(i)} \\ y_{13}^{(i)} & y_{13}^{(i)} & y_{33}^{(i)} \end{bmatrix} \begin{bmatrix} F_1^{(i)} \\ F_2^{(i)} \\ V_3^{(i)} \end{bmatrix} \quad (2)$$

while the interconnection laws are (for mechanical ports, and $i > 1$)

$$\left. \begin{aligned} F_1^{(i)} &= F_2^{(i-1)} \\ V_1^{(i)} &= -V_2^{(i-1)} \end{aligned} \right\} \quad (3)$$

where

$$\left. \begin{aligned} F_1^{(1)} &= F_{in} = F_{left\ boundary}, & V_1^{(1)} &= V_{in} \\ F_2^{(N)} &= F_{out} = F_{right\ boundary}, & V_2^{(N)} &= V_{out} \end{aligned} \right\} \quad (4)$$

Also for electrical ports with $i \geq 1$, the interconnection laws are

$$\begin{aligned} V_3^{(i)} &= (-1)^{i-1} V_3 \\ I_3 &= \sum_{i=1}^N (-1)^{i-1} I_3^{(i)} \end{aligned} \quad (5)$$

In order to find the interconnected network 3-port description we convert to a reverse - chain matrix type of description. From equation (2) it follows that

$$\begin{bmatrix} F_2^{(i)} \\ -V_2^{(i)} \end{bmatrix} = Y_R^{(i)} \begin{bmatrix} F_1^{(i)} \\ V_1^{(i)} \end{bmatrix} + A^{(i)} V_3 \quad (6)$$

technique by using
n. The starting
ion of the SAW
e matrix of a full
nterdigital period
ce matrix for an
ions in cascade
and admittance
as expressed in

NETWORK
are recorded here
ports (of variables

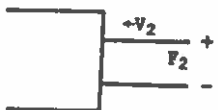


Diagram of a three-port

y_{13} by
nical ports and in
orts as shown in

where

$$\psi_R^{(1)} = -\frac{1}{y_{12}^{(1)}} \begin{bmatrix} y_{11}^{(1)} & -1 \\ (y_{12}^{(1)})^2 - (y_{11}^{(1)})^2 & y_{11}^{(1)} \end{bmatrix} \quad (7)$$

$$A^{(1)} = \frac{(-1)^1 y_{13}^{(1)}}{y_{12}^{(1)}} \begin{bmatrix} 1 \\ y_{12}^{(1)} - y_{11}^{(1)} \end{bmatrix} \quad (8)$$

and

$$I_3^{(1)} = -\frac{y_{13}^{(1)}}{y_{12}^{(1)}} [y_{11}^{(1)} - y_{12}^{(1)}, 1] \begin{bmatrix} F_1^{(1)} \\ v_1^{(1)} \end{bmatrix} + \frac{(-1)^{1-1}}{y_{12}^{(1)}} [y_{33}^{(1)} y_{12}^{(1)} - (y_{13}^{(1)})^2] v_3 \quad (9)$$

Equation (6) is the key for this analysis. Using the mechanical interconnections, we obtain

$$\begin{bmatrix} F_2^{(1)} \\ -v_2^{(1)} \end{bmatrix} = \psi_R^{(1)} \begin{bmatrix} F_{1n} \\ v_{1n} \end{bmatrix} + A^{(1)} v_3 \quad (10)$$

$$\begin{bmatrix} F_2^{(2)} \\ -v_2^{(2)} \end{bmatrix} = \psi_R^{(2)} \begin{bmatrix} F_1^{(2)} \\ v_1^{(2)} \end{bmatrix} + A^{(2)} v_3 \quad (11)$$

or

$$\begin{bmatrix} F_2^{(2)} \\ -v_2^{(2)} \end{bmatrix} = \psi_R^{(2)} \psi_R^{(1)} \begin{bmatrix} F_{1n} \\ v_{1n} \end{bmatrix} + (\psi_R^{(2)} A^{(1)} + A^{(2)}) v_3 \quad (12)$$

In general for N Sections

$$\begin{bmatrix} F_{out} \\ -v_{out} \end{bmatrix} = \begin{bmatrix} F_2^{(N)} \\ -v_2^{(N)} \end{bmatrix} = \psi_R^{(N)} \begin{bmatrix} F_1^{(N)} \\ -v_1^{(N)} \end{bmatrix} + A^{(N)} v_3$$

$$= \psi_R^{(N)} \psi_R^{(N-1)} \dots$$

$$\dots \psi_R^{(3)} A^{(2)} + \psi_R^{(N)}$$

Thus for the entire

where

The notation \prod_R \equiv identity. If all (16) that

$$P_R \equiv (\psi_R^{(1)})^{N-1}$$

and

$$A = \sum_{i=1}^N (\psi_R^{(1)})^{N-i} A^{(i)}$$

where $N-i \equiv k$. Cont. section

$$\psi_R = -\frac{1}{y_{12}}$$

so that

$$\psi_R^2 = \frac{1}{y_{12}^2}$$

$$\begin{aligned}
 (7) \quad & - \psi_R^{(N)} \psi_R^{(N-1)} \dots \psi_R^{(2)} \psi_R^{(1)} \begin{bmatrix} v_{in} \\ v_{in} \end{bmatrix} + (A^{(N)} + \psi_R^{(N)} A^{(N-1)} + \dots + \psi_R^{(N)} \psi_R^{(N-1)} \dots \\
 & \dots \psi_R^{(3)} A^{(2)} + \psi_R^{(N)} \psi_R^{(N-1)} \dots \psi_R^{(3)} \psi_R^{(2)} A^{(1)}) \quad (13)
 \end{aligned}$$

Thus for the entire cascaded system

$$(8) \quad \begin{bmatrix} F_{out} \\ -v_{out} \end{bmatrix} = P_R \begin{bmatrix} F_{in} \\ v_{in} \end{bmatrix} + A v_3 \quad (14)$$

where

$$P_R = \prod_{i=1}^N \psi_R^{(i)} \quad (15)$$

$$A = \sum_{i=1}^N \left(\prod_{k=1}^{N-1} \psi_R^{(k+1)} \right) A^{(i)} \quad (16)$$

(10) The notation \prod_R means multiply lower indexed terms on the right and $\prod_{k=N}^{N-1}$ = identity. If all sections are identical then it follows from (7), (15) and (16) that

$$(11) \quad P_R = (\psi_R^{(1)})^N = \frac{(-1)^N}{y_{12}^2} \begin{bmatrix} y_{11} & -1 \\ y_{12}^2 - y_{11}^2 & y_{11} \end{bmatrix} \quad (17)$$

and

$$(12) \quad A = \sum_{i=1}^N (\psi_R^{(1)})^{N-1} (-1)^{i-1} A^{(i)} = \sum_{k=0}^{N-1} (\psi_R^{(1)})^k (-1)^{N-k} \frac{y_{13}}{y_{12}} \begin{bmatrix} 1 \\ y_{12} - y_{11} \end{bmatrix} \quad (18)$$

where $N-1 \equiv k$. Continuing with the identical section case we have for each section

$$\psi_R = -\frac{1}{y_{12}} \begin{bmatrix} y_{11} & -1 \\ y_{12}^2 - y_{11}^2 & y_{11} \end{bmatrix}, \quad A = -\frac{y_{13}}{y_{12}} \begin{bmatrix} 1 \\ y_{12} - y_{11} \end{bmatrix} \quad (19)$$

so that

$$\psi_R^2 = \frac{1}{y_{12}^2} \begin{bmatrix} y_{11} & -1 \\ y_{12}^2 - y_{11}^2 & y_{11} \end{bmatrix} \begin{bmatrix} y_{11} & -1 \\ y_{12}^2 - y_{11}^2 & y_{11} \end{bmatrix} \quad (20)$$

$-(y_{13}^{(1)})^2] v_3$, (9)

Interconnec-

or

$$\psi_R^2 = \frac{1}{y_{12}^2} \begin{bmatrix} 2y_{11}^2 - y_{12}^2 & -2y_{11} \\ 2y_{11}(y_{12}^2 - y_{11}^2) & 2y_{11}^2 - y_{12}^2 \end{bmatrix} \quad (21)$$

Thus it leads to the recursion relation

$$\psi_R^2 = -\psi_R^0 - \frac{2y_{11}}{y_{12}} \psi_R^1 \quad (22)$$

where

$$\psi_R^0 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \quad (23)$$

Relation (22) can be used to reduce the higher powers of ψ_R to the first.

We next proceed to find the admittance matrix of the overall 3-port network.

Explicit expansion of (14) yields

$$\begin{bmatrix} v_{in} \\ v_{out} \end{bmatrix} = \begin{bmatrix} -P_{R12}^{-1} P_{R11} & P_{R12}^{-1} & -P_{R12}^{-1} A_{11} \\ -P_{R12}^{-1} + P_{R22} P_{R12}^{-1} & P_{R11} & -P_{R22} P_{R12}^{-1} A_{11} - A_{21} \end{bmatrix} \begin{bmatrix} F_{in} \\ F_{out} \\ v_3 \end{bmatrix} \quad (24)$$

where P_{R11} , P_{R12} , P_{R22} etc. represent the corresponding elements of the final matrix P_R for the entire cascaded network system.

Following the use of induction, an explicit evaluation of the matrix elements in (24) gives the following results for all $N = 1, 2, 3, \dots$

$$Y_{13} = y_{13} \quad (25)$$

$$Y_{23} = (-1)^{N-1} y_{13} \quad (26)$$

Y_{ij} are the matrix elements of the final admittance matrix Y for the N cascaded sections, and y_{ij} are the matrix elements for the single electrode section. It remains to calculate I_3 to complete the admittance matrix. We have from (5)

$$I_3 = \sum_{i=1}^N (-1)^{i-1} I_3^{(i)} = \sum_{i=1}^N (-1)^{i-1} y_{13}^{(i)} [F_1^{(i)} + F_2^{(i)}] + \sum_{i=1}^N y_{33}^{(i)} v_3 \quad (27)$$

where

$$F_1^{(i)} = F_2^{(i-1)} \quad (28)$$

$$\begin{bmatrix} F_2^{(i)} \\ -v_2^{(i)} \end{bmatrix} = \prod_{j=1}^i \psi_R^{(j)} \begin{bmatrix} F_{in} \\ v_{in} \end{bmatrix} + \sum_{j=1}^i \left(\prod_{k=j}^{i-1} \psi_R^{(k+1)} \right) A^{(j)} v_3 \quad (29)$$

and

$$\begin{bmatrix} F_1^{(i)} \\ -v_2^{(i-1)} \end{bmatrix} = \prod_{j=1}^{i-1} \psi_R^{(j)}$$

for which the (1,1) substituted.

For identical sec

$$\begin{bmatrix} F_2^{(i)} \\ -v_2^{(i)} \end{bmatrix} = \frac{(-1)^{i-1}}{y_{12}^{i-1}}$$

$$+ \left(\sum_{k=0}^{i-2} (\psi_R^{(1)})^k \right)$$

$$\begin{bmatrix} F_1^{(i)} \\ -v_2^{(i-1)} \end{bmatrix} = \frac{(-1)^{i-1}}{y_{12}^{i-1}}$$

$$+ \sum_{k=0}^{i-2} (\psi_R^{(1)})^k$$

After some algebraic $[F_1^{(i)} + F_2^{(i)}]$. 0 (27) we get

$$I_3 = \sum_{i=1}^N y_{33} v_3$$

$$-(y_{12} - y_{11}) [y_{12} + y_{11} - (y_{12} - y_{11}) (\psi_R^{(1)})^{i-1}]$$

From this result for and $Y_{32} = Y_{23}$, can reciprocity relation. Therefore, the final equal section case

$$\begin{bmatrix} V_1^{(1)} \\ -V_2^{(1-1)} \end{bmatrix} = \prod_{j=1}^{i-1} \psi_R^{(j)} \begin{bmatrix} V_{in} \\ V_{in} \end{bmatrix} + \sum_{j=1}^{i-1} \left(\prod_{k=j}^{i-2} \psi_R^{(k+1)} \right) A^{(j)} v_3 \tag{21}$$

for which the (1,1) entries are to be summed and the expression for V_{in} substituted.

(22)

For identical sections, it turns out that

(23)

$$\begin{bmatrix} V_2^{(1)} \\ -V_2^{(1)} \end{bmatrix} = \frac{(-1)^{i-1}}{y_{12}^{i-1}} \begin{bmatrix} y_{11} & -1 \\ y_{12}^2 - y_{11}^2 & y_{11} \end{bmatrix}^{i-1} \frac{(-1)}{y_{12}} \begin{bmatrix} y_{11} & -1 \\ y_{12}^2 - y_{11}^2 & y_{11} \end{bmatrix} \begin{bmatrix} V_{in} \\ V_{in} \end{bmatrix} + \left(\sum_{k=0}^{i-2} (\psi_R^{(1)})^k (-1)^{i-k} + (\psi_R^{(1)})^{i-1} (-1)^{i-(i-1)} \right) \frac{y_{13}}{y_{12}} \begin{bmatrix} 1 \\ y_{12} - y_{11} \end{bmatrix} v_3 \tag{31}$$

to the first. all 3-port network.

$$\begin{bmatrix} V_{in} \\ V_{out} \\ v_3 \end{bmatrix} \tag{24}$$

elements of the the matrix 2, 3,

(25)

(26)

$$\begin{bmatrix} V_1^{(1)} \\ -V_2^{(1-1)} \end{bmatrix} = \frac{(-1)^{i-1}}{y_{12}^{i-1}} \begin{bmatrix} y_{11} & -1 \\ y_{12}^2 - y_{11}^2 & y_{11} \end{bmatrix}^{i-1} \begin{bmatrix} V_{in} \\ V_{in} \end{bmatrix} + \sum_{k=0}^{i-2} (\psi_R^{(1)})^k (-1)^{i-1-k} \frac{y_{13}}{y_{12}} \begin{bmatrix} 1 \\ y_{12} - y_{11} \end{bmatrix} v_3 \tag{32}$$

After some algebraic manipulation of (31) and (32), we find an expression for $[V_1^{(1)} + V_2^{(1)}]$. On substitution of this expression for $[V_1^{(1)} + V_2^{(1)}]$ into (27) we get

for the N cascaded electrode section. It We have from (5)

$$y_{33}^{(1)} v_3 \tag{27}$$

(28)

(29)

$$\begin{aligned}
 I_3 = & \sum_{i=1}^N y_{33} v_3 + \sum_{i=1}^N (-1)^{i-1} \frac{y_{13}}{y_{12}} [(y_{12} - y_{11} - P_{R12}^{-1} P_{R11}) (\psi_R^{(1)})^{i-1}]_{11} \\
 & - (y_{12} - y_{11}) [(y_{12} + y_{11} - (-P_{R12}^{-1} P_{R11})) (\psi_R^{(1)})^{i-1}]_{12} v_{in} + ((\psi_R^{(1)})^{i-1})_{11} \\
 & + (y_{12} - y_{11}) (\psi_R^{(1)})^{i-1} P_{R12}^{-1} v_{out} \tag{33}
 \end{aligned}$$

From this result for $N = 1, 2, 3$ the expected reciprocity results, $Y_{31} = Y_{13}$ and $Y_{32} = Y_{23}$, can be easily verified. By induction it follows that the reciprocity relations are valid for an arbitrary N .

Therefore, the final simplified form of the admittance matrix Y for N equal section case is given by

$$Y = \begin{bmatrix} -P_{R12}^{-1} P_{R11} & P_{R12}^{-1} & y_{13} \\ P_{R12}^{-1} & -P_{R12}^{-1} P_{R11} & (-1)^{N-1} y_{13} \\ y_{13} & (-1)^{N-1} y_{13} & N y_{33} \end{bmatrix} \quad (34)$$

This is the most general form of the admittance matrix for the "N cascaded" network shown in Fig. 2. The combined network structure consequently may be defined by a single section as shown in Fig. 3.

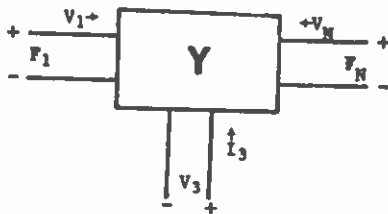


Figure 3. Combined Network Structure

3. EQUIVALENT CIRCUIT OF THE IN-LINE MODEL AND THE ADMITTANCE MATRIX

The development of the interdigital transducer transfer function begins with the one dimensional in-line model for a single electrode section; this is shown for the sinusoidal steady state in Fig. 4. The model of a single electrode section has three ports, corresponding to the two symmetric acoustic ports and the electric port. The stresses and particle velocities at the acoustic ports are represented by equivalent voltages and currents, respectively, to which they are numerically equal. A convenient method for describing the relationship between these currents and voltages, $V_j = y_{jk} F_k$, is in terms of the admittance matrix as for equation (1). The y_{jk} 's for the nth section are given by [5]

$$y_{11} = \frac{1}{jZ_o} \frac{(\cos \theta_n - \nu \sin \theta_n)}{\sin \theta_n + 2\nu(\cos \theta_n - 1)} = y_{22} \quad (35)$$

$$y_{12} = \frac{1}{Z_o} \frac{(1 - \nu \sin \theta_n)}{\sin \theta_n + 2\nu(\cos \theta_n - 1)} = y_{21} \quad (36)$$

$$y_{13} = \frac{j\nu}{Z_o} \frac{(\cos \theta_n - 1)}{\sin \theta_n + 2\nu(\cos \theta_n - 1)} = y_{23} = y_{32} = y_{31} \quad (37)$$

$$y_{33} = j\omega C_o \left[1 + \dots \right]$$

where



Figure 4. Equival

$$y_{33} = j\omega C_o \left[1 + \frac{2K^2}{\theta_n} \cdot \frac{1 - \cos \theta_n}{\sin \theta_n + 2\gamma(\cos \theta_n - 1)} \right] \quad (38)$$

where

$\theta_n = \frac{l_n}{V} \cdot 2\pi f$, is the acoustic transit angle of the n^{th} section; f = frequency.

$\phi = \left(\frac{K^2 C_o Z_o V}{l_n} \right)^{1/2}$ is the transformer turns ratio.

$\gamma = \frac{K^2}{\theta_n}$

$C_o = \frac{C}{\phi^2}$, C_o = capacitance of the n^{th} section

$\omega = 2\pi f$

Z_o = acoustic characteristic impedance

K^2 = piezo-electric coupling coefficient

V = velocity of acoustic propagation

l_n = length of n^{th} section

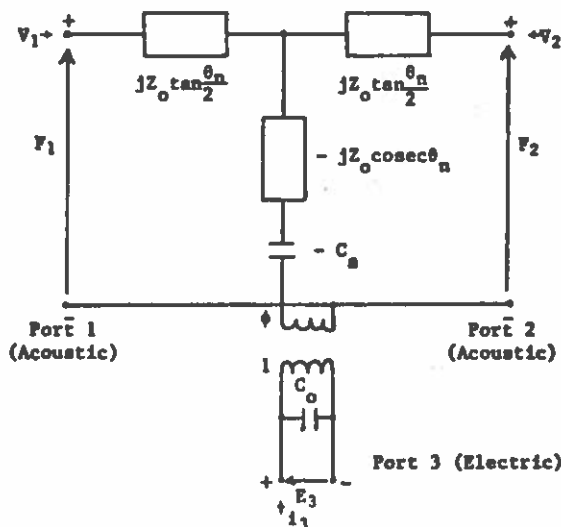


Figure 4. Equivalent circuit of the in-line model for one-electrode section.

3
33
he "N cascaded"
sequently may be

ion begins with the
this is shown for
electrode section
ports and the
acoustic ports are
, to which they
he relationship
of the admittance
e given by [5]

(35)

(36)

(37)

4. SCATTERING MATRIX FOR N-SECTION IN-LINE MODEL

It is well known [6,7] that the scattering matrix of a network characterized by its admittance matrix Y is

$$S = \underline{I}_3 - 2Y (\underline{I}_3 + Y)^{-1} \tag{39}$$

where \underline{I}_3 is the 3x3 identity matrix. This equation (39) will be used to find the scattering matrix for the N-section in-line model.

Application of a unitary transformation for diagonalization of $\Phi_R^{(1)}$ given by (7) gives the product matrix P_R for the N-section in-line model as

$$P_R = (\Phi_R^{(1)})^N = \sin N\delta \begin{bmatrix} \cot N\delta & -jZ_0 \sin \delta (1 - \gamma \sin \delta) (\sin \delta)^{-1} \\ -\frac{1}{Z_0} \frac{\sin \delta}{\sin \delta (1 - \gamma \sin \delta)} & \cot N\delta \end{bmatrix} \tag{40}$$

where δ is given by the expression

$$\sin \delta = (\sin \theta + 2\gamma \cos \theta - 2\gamma)^{1/2} (1 - \gamma \sin \theta)^{-1/2} (\sin \theta)^{1/2} \tag{41}$$

Using (34), (37), (38) and (40), the admittance matrix Y for the N-section in-line model is obtained in the final form

$$Y = \begin{bmatrix} -\frac{1}{Z_0} \frac{\cot N\delta}{(1-2\gamma \tan \frac{\theta}{2})^{1/2}} & \frac{1}{Z_0} \frac{\operatorname{cosec} N\delta}{(1-2\gamma \tan \frac{\theta}{2})^{1/2}} & -\frac{1}{Z_0} \left\{ \frac{\tan \frac{\theta}{2}}{(1-2\gamma \tan \frac{\theta}{2})} \right\} \\ \frac{1}{Z_0} \frac{\operatorname{cosec} N\delta}{(1-2\gamma \tan \frac{\theta}{2})^{1/2}} & -\frac{1}{Z_0} \frac{\cot N\delta}{(1-2\gamma \tan \frac{\theta}{2})^{1/2}} & (-1)^N \frac{1}{Z_0} \left\{ \frac{\tan \frac{\theta}{2}}{(1-2\gamma \tan \frac{\theta}{2})} \right\} \\ -\frac{1}{Z_0} \left\{ \frac{\tan \frac{\theta}{2}}{(1-2\gamma \tan \frac{\theta}{2})} \right\} & (-1)^N \frac{1}{Z_0} \left\{ \frac{\tan \frac{\theta}{2}}{(1-2\gamma \tan \frac{\theta}{2})} \right\} & j\omega C_T \left(1 + \frac{2K^2}{\theta} \frac{\tan \frac{\theta}{2}}{(1-2\gamma \tan \frac{\theta}{2})} \right) \end{bmatrix} \tag{42}$$

where $C_T = NC_0$.

Finally, direct substitution of Y from (42) into (39) yields the explicit form of the scattering matrix for the N-section in-line model. A straightforward but lengthy calculations [6] gives the explicit and final form of the scattering matrix elements in terms of known parameters as

$$S_{11} = S_{22} = \frac{1}{|M|} \left[\left(1 - \frac{1}{n^2 Z_0^2} \right) + j \left\{ \omega C_T \left(1 - \frac{1}{n^2 Z_0^2} \right) \cdot \left(1 + \frac{2K^2}{n^2} \tan \frac{\theta}{2} \right) + \frac{2\phi^2}{n^2 Z_0^3} \tan^2 \left(\frac{\theta}{2} \right) (\cot N\delta + (-1)^{N+1} \operatorname{cosec} N\delta) \right\} \right] \tag{43}$$

$$S_{12} = S_{21}$$

$$S_{13} = S_{31}$$

$$S_{23} = S_{32}$$

and

$$S_{33} = \frac{1}{|M|} \left[\left(1 + \frac{1}{n^2 Z_0^2} \right) \right.$$

$$\left. + j \left\{ \frac{2\phi^2}{n^2} \right. \right.$$

$$\left. - \omega C_T \left(1 + \frac{1}{n^2 Z_0^2} \right) \right]$$

where

$$|M| = \left[\left(1 + \frac{1}{n^2 Z_0^2} \right) \right.$$

$$\left. + j \left\{ \omega C_T \left(1 + \frac{1}{n^2 Z_0^2} \right) \right. \right.$$

$$\left. - \frac{2\phi^2}{n^2 Z_0^3} \tan^2 \left(\frac{\theta}{2} \right) \right]$$

and

Therefore, the scattering matrix elements for reciprocal circuits are rather bulky, the

work characterized by

(39)

will be used to find

ion of $\psi_R^{(1)}$ given by
model as

$$\left. \begin{aligned} & \sin \theta (\sin \theta)^{-1} \\ & \text{at } N\delta \end{aligned} \right\} (40)$$

$$i^{-1} (\sin \theta)^{1/2} \quad (41)$$

for the N-section

$$\left. \begin{aligned} & \frac{\tan \frac{\theta}{2}}{(1-2\gamma \tan \frac{\theta}{2})} \\ & \frac{\tan \frac{\theta}{2}}{(1-2\gamma \tan \frac{\theta}{2})} \\ & 1 + \frac{2K^2}{\theta} \frac{\tan \frac{\theta}{2}}{(1-2\gamma \tan \frac{\theta}{2})} \end{aligned} \right\} (42)$$

gives the explicit form
A straightforward
form of the scattering

$$\left. \begin{aligned} & \frac{1}{Z_0^2} \\ & N\delta + (-1)^{N+1} \operatorname{cosec} N\delta \end{aligned} \right\} (43)$$

$$S_{12} = S_{21} = \frac{2}{|M|} \left[\left\{ \frac{\omega C_T}{n Z_0} \left(1 + \frac{2K^2}{n^2 \theta} \tan \frac{\theta}{2} \right) \operatorname{cosec} N\delta \right. \right. \\ \left. \left. + (-1)^N \frac{\phi^2}{n^4 Z_0^2} \tan^2 \left(\frac{\theta}{2} \right) \right\} - \frac{1}{n Z_0} \operatorname{cosec} N\delta \right] \quad (44)$$

$$S_{13} = S_{31} = \frac{2}{|M|} \left[\frac{\phi}{n^3 Z_0^2} \tan \frac{\theta}{2} \left\{ \cot N\delta + (-1)^{N+1} \operatorname{cosec} N\delta \right\} \right. \\ \left. + j \frac{\phi^2}{n^2 Z_0} \tan \frac{\theta}{2} \right] \quad (45)$$

$$S_{23} = S_{32} = (-1)^{N-1} S_{13} \quad (46)$$

and

$$S_{33} = \frac{1}{|M|} \left[\left\{ \left(1 + \frac{1}{n^2 Z_0^2} \right) - \frac{2\phi^2}{n^4 Z_0^2} \tan^2 \left(\frac{\theta}{2} \right) - \frac{2\omega C_T}{n Z_0} \cot N\delta \left(1 + \frac{2K^2}{n^2 \theta} \tan \frac{\theta}{2} \right) \right\} \right. \\ \left. + j \left\{ \frac{2\phi^2}{n^5 Z_0^3} \tan^2 \left(\frac{\theta}{2} \right) (\cot N\delta + (-1)^{N+1} \operatorname{cosec} N\delta) \right. \right. \\ \left. \left. - \omega C_T \left(1 + \frac{1}{n^2 Z_0^2} \right) \left(1 + \frac{2K^2}{n^2 \theta} \tan \frac{\theta}{2} \right) - \frac{2}{n Z_0} \cot N\delta \right\} \right] \quad (47)$$

where

$$|M| = \left[\left\{ \left(1 + \frac{1}{n^2 Z_0^2} \right) + \frac{2\phi^2}{n^4 Z_0^2} \tan^2 \left(\frac{\theta}{2} \right) + \frac{2\omega C_T}{n Z_0} \cot N\delta \left(1 + \frac{2K^2}{n^2 \theta} \tan \frac{\theta}{2} \right) \right\} \right. \\ \left. + j \left\{ \omega C_T \left(1 + \frac{1}{n^2 Z_0^2} \right) \left(1 + \frac{2K^2}{n^2 \theta} \tan \frac{\theta}{2} \right) - \frac{2}{n Z_0} \cot N\delta \right. \right. \\ \left. \left. - \frac{2\phi^2}{n^5 Z_0^3} \tan^2 \left(\frac{\theta}{2} \right) (\cot N\delta + (-1)^{N+1} \operatorname{cosec} N\delta) \right\} \right] \quad (48)$$

$$\text{and} \quad \eta = (1 - 2\gamma \tan \frac{\theta}{2})^{1/2} \quad (49)$$

Therefore, the scattering matrix for the N-section in-line model reflects the circuits reciprocity in its symmetric form. Although these expressions are rather bulky, they should be useful for computer aided analysis.

5. CONCLUSION

Here we have given, for the sinusoidal steady state behavior, explicit formulas for the scattering matrix entries of a SAW interdigital transducer based upon the in-line model of a one-electrode section, these formulas appearing in equations (43) - (49). In the form presented, they are suitable for computer aided design of SAW transducers while their existence paves the way for more extensive design formalisms.

6. REFERENCES

- [1] H. Mathews, ed. *Surface Wave Filters*, see ch. 2, R.F. Milson, M. Redwood and N.H.C. Reilly, "The Interdigital Transducer", New York: John Wiley, 1977.
- [2] A.A. Oliner, ed. *Acoustic Surface Waves*, see ch. 3, H.M. Gerard, "Principles of Surface Wave Filter Design", Berlin: Springer-Verlag, 1978.
- [3] W.R. Smith, H.M. Gerard, J.H. Collins, T.M. Reeder, and H.J. Shaw, "Analysis of Interdigital Surface Wave Transducers by Use of an Equivalent Circuit Model", *IEEE Trans. Microwave Theory and Techniques*, Vol. MTT-17, pp. 856-864, November 1969.
- [4] H.M. Gerard, "Acoustic Scattering Parameters of the Electrically Loaded Interdigital Surface Wave Transducer", *IEEE Trans. Microwave Theory and Techniques*, Vol. MTT-17, pp. 1045-1046, November 1969.
- [5] M. Hribsek and R.W. Newcomb, "The Scattering Matrix of a SAW section", *Proceedings of the Fourth International Symposium on Network Theory*, Ljubljana, pp. 293-298, September 1979.
- [6] N.C. Debnath, R.C. Ajmera, and R.W. Newcomb, "Scattering Matrix for N-section Cross-field Model of SAW Transducers", *Proceedings of the 12th Southeastern Symposium on System Theory*, Virginia, May 1980.
- [7] R.W. Newcomb, *Linear Multiport Synthesis*, p. 82, New York: McGraw Hill, 1966.

Address of the authors:

N.C. Debnath and R.C. Ajmera, Department of Physics, East Carolina University Greenville, North Carolina 27834, U.S.A.

R.W. Newcomb, Department of Electrical Engineering, University of Maryland, College Park, Maryland 20742, U.S.A.

Miloslav Košek
Technical College

SURFACE AC
BY

A superposition of a bandpass filter and a bandstop filter apodized transducer block with a transfer function suitable one. This is a filter with both pass and stop phase characteristics determined by means of the scattering matrix of the transducer circuit. The exactness of the design is determined by the tolerance

The main area of application of the filter is for the television receiver. The asymmetric amplitude response is both specified by the pass and stop phase characteristics.

The fundamental method of the design is the method of the equivalent circuit. The transfer function suitable for the IF filter is determined by the tolerance field. The design is very severe in the pass band and less severe in the stop band. The thesis of the filter [4] or its modification.

In this paper we present a suitable configuration for the results of its application with

ECCTD'80

**1980 EUROPEAN CONFERENCE
ON CIRCUIT THEORY AND DESIGN**

Proceedings

Vol. 2

Warsaw, Poland. September 2-5, 1980